

An H^1 -Galerkin mixed method for the coupled Burgers equation

Xianbiao Jia, Hong Li, Yang Liu, Zhichao Fang

Abstract—In this paper, an H^1 -Galerkin mixed finite element method is discussed for the coupled Burgers equations. The optimal error estimates of the semi-discrete and fully discrete schemes of the coupled Burgers equation are derived.

Keywords—The coupled Burgers equation; H^1 -Galerkin mixed finite element method; Backward Euler's method; Optimal error estimates.

I. INTRODUCTION

WITH the research and development of the mixed finite element methods and H^1 -Galerkin method, Pani [2] (in 1998) proposed a new mixed finite element method called H^1 -Galerkin mixed finite element procedure which is applied to a mixed system in u and its flux q . The approximating finite element spaces V_h and W_h are allowed to be of differing polynomial degrees. Hence, estimations have been obtained which distinguish the better approximation properties of V_h and W_h . Compared to standard mixed methods, the proposed one is not subject to LBB consistency condition. Although we require extra regularity on the solution, a better order of convergence for the flux in L^2 -norm is obtained. From then on, the method was applied to the evolution integro-differential equation[3],[4],[5],[6], hyperbolic problems[9],[10],[11],[14],[16], fourth-order parabolic equation[15], Sobolev equation[7],[8], Schrodinger equation[13] and nonlinear evolution equations[17],[18],[12] and so on. In this paper, we propose H^1 -Galerkin mixed finite element scheme for the following coupled Burgers equation[1]

$$\begin{cases} u_t - u_{xx} - 2uu_x + (uv)_x = f(x, t), (x, t) \in \Omega \times J, \\ v_t - v_{xx} - 2vv_x + (uv)_x = g(x, t), (x, t) \in \Omega \times J, \\ u(x, t) = 0, v(x, t) = 0, (x, t) \in \partial\Omega \times \bar{J}, \\ u(x, 0) = u_0(x), v(x, 0) = v_0(x), x \in \Omega, \end{cases} \quad (1)$$

where $\Omega = [0, 1] \subset \mathbb{R}^1$ with Lipschitz continuous boundary $\partial\Omega$, $J = (0, T]$ is the time interval with $0 < T < \infty$, $f(x, t), g(x, t)$ are two functions.

II. H^1 -GALERKIN MIXED FINITE ELEMENT METHOD

Denote the natural inner product on $L^2(I)$ as $(\cdot, \cdot)$. Let $H_0^1 = \{z \in H^1(I) \mid z(0) = z(1) = 0\}$. Further, we call the classical Sobolev spaces $W^{m,p}(I)$, $1 \leq p \leq \infty$ as $W^{m,p}$

School of Mathematical Sciences, Inner Mongolia University, Hohhot 010021, China. Correspondence to: E-mail: smslh@imu.edu.cn (H. Li); math-liuyang@yahoo.cn (Y. Liu).

Manuscript received April 22, 2012.

with norm $\|\cdot\|_{m,p}$. When $p = 2$, we simply write $W^{m,p}$ as H^m with norm $\|\cdot\|_m$.

With $q = u_x, \sigma = v_x$ we reformulate the formulation (1) as the first-order system:

$$\begin{cases} u_x = q, v_x = \sigma, \\ u_t - q_x - 2uq + qv + u\sigma = f(x, t), \\ v_t - \sigma_x - 2v\sigma + \sigma u + qv = g(x, t), \end{cases} \quad (2)$$

To derive the H^1 -Galerkin mixed finite element method, we consider the following weak formulation of (1): find $\{u, v, q, \sigma\} : [0, T] \rightarrow H_0^1 \times H^1$ satisfying:

$$\begin{cases} (u_x, \chi_x) = (q, \chi_x), \forall \chi \in H_0^1, (a) \\ (v_x, w_x) = (\sigma, w_x), \forall w \in H_0^1, (b) \\ (q_t, \phi) + (q_x, \phi_x) + 2(uq, \phi_x) - (qv, \phi_x) - (u\sigma, \phi_x) \\ = -(f, \phi_x), \forall \phi \in H^1, (c) \\ (\sigma_t, \psi) + (\sigma_x, \psi_x) + 2(v\sigma, \psi_x) - (\sigma u, \psi_x) - (qv, \psi_x) \\ = -(g, \psi_x), \forall \psi \in H^1, (d) \end{cases} \quad (3)$$

For (3c,d), we have used integration by parts, and the *Dirichlet* boundary conditions $u_t(0, t) = u_t(1, t) = 0, v_t(0, t) = v_t(1, t) = 0$.

Let V_h and W_h be finite dimensional subspaces of H_0^1 and H^1 , respectively, with the following approximation properties: for $1 \leq p \leq \infty$ and k, r positive integers

$$\begin{aligned} & \inf_{v_h \in V_h} \{ \|v - v_h\|_{L^p} + h \|v - v_h\|_{W^{1,p}} \} \\ & \leq Ch^{k+1} \|v\|_{W^{k+1,p}}, v \in H_0^1 \cap W^{k+1,p}, \\ & \inf_{w_h \in W_h} \{ \|w - w_h\|_{L^p} + h \|w - w_h\|_{W^{1,p}} \} \\ & \leq Ch^{r+1} \|w\|_{W^{r+1,p}}, w \in W^{r+1,p}. \end{aligned}$$

The semidiscrete H^1 -Galerkin mixed finite element for (3) consists in determining $\{u^h, v^h, q^h, \sigma^h\} : [0, T] \rightarrow V_h \times W_h$ such that:

$$\begin{cases} (u_x^h, \chi_x^h) = (q^h, \chi_x^h), \forall \chi^h \in V_h, (a) \\ (v_x^h, w_x^h) = (\sigma^h, w_x^h), \forall w^h \in W_h, (b) \\ (q_t^h, \phi^h) + (q_x^h, \phi_x^h) + 2(u^h q^h, \phi_x^h) - (q^h v^h, \phi_x^h) \\ - (u^h \sigma^h, \phi_x^h) = -(f^h, \phi_x^h), \forall \phi^h \in W_h, (c) \\ (\sigma_t^h, \psi^h) + (\sigma_x^h, \psi_x^h) + 2(v^h \sigma^h, \psi_x^h) - (\sigma^h u^h, \psi_x^h) \\ - (q^h v^h, \psi_x^h) = -(g^h, \psi_x^h), \forall \psi^h \in W_h, (d) \end{cases} \quad (4)$$

For use in the error analysis, we define the elliptic projection $\tilde{u}^h, \tilde{v}^h \in V_h$ by

$$(u_x - \tilde{u}_x^h, \chi_x^h) = 0, (v_x - \tilde{v}_x^h, \omega_x^h) = 0, \chi^h, \omega^h \in V_h. \quad (5)$$

Further, we also define a Ritz projection $\tilde{q}^h, \tilde{\sigma}^h \in W_h$ of q, σ as the solution of

$$A(q - \tilde{q}^h, \phi^h) = 0, A(\sigma - \tilde{\sigma}^h, \psi^h) = 0, \phi^h, \psi^h \in W_h. \quad (6)$$

where $A(z, w) = (z_x, w_x) + \lambda(z, w)$. Here λ is chosen appropriately so that A is H^1 -coercive, i.e.,

$$A(w, w) \geq \mu_0 \|w\|_1^2, w \in H^1$$

where μ_0 is a positive constant. Moreover, it is not hard to check that $A(\cdot, \cdot)$ is bounded.

With $\eta = u - \tilde{u}^h, \tau = v - \tilde{v}^h, \rho = q - \tilde{q}^h, \delta = \sigma - \tilde{\sigma}^h$, the following estimates are well known [19]: for $j = 0, 1$

$$\|\eta\|_j + \|\eta_t\|_j \leq Ch^{k+1-j} [\|u\|_{k+1} + \|u_t\|_{k+1}], \|\tilde{u}^h\|_{0,\infty} \leq C(u) \quad (7)$$

$$\|\tau\|_j + \|\tau_t\|_j \leq Ch^{k+1-j} [\|v\|_{k+1} + \|v_t\|_{k+1}], \|\tilde{v}^h\|_{0,\infty} \leq C(v) \quad (8)$$

$$\|\rho\|_j \leq Ch^{r+1-j} \|q\|_{r+1}, \|\rho_t\|_j \leq Ch^{r+1-j} \|q_t\|_{r+1} \quad (9)$$

$$\|\delta\|_j \leq Ch^{r+1-j} \|\sigma\|_{r+1}, \|\delta_t\|_j \leq Ch^{r+1-j} \|\sigma_t\|_{r+1} \quad (10)$$

III. ERROR ESTIAMTES FOR SEMI-DISCRETE SCHEME

For a priori error estimates, we decompose the errors as $u - u^h = u - \tilde{u}^h + \tilde{u}^h - u^h = \eta + \varsigma$; $v - v^h = v - \tilde{v}^h + \tilde{v}^h - v^h = \tau + \theta$; $q - q^h = q - \tilde{q}^h + \tilde{q}^h - q^h = \rho + \xi$; $\sigma - \sigma^h = \sigma - \tilde{\sigma}^h + \tilde{\sigma}^h - \sigma^h = \delta + \gamma$. From (3)-(6), we then obtain

$$\begin{cases} (\varsigma_x, \chi_x^h) = (\rho, \chi_x^h) + (\xi, \chi_x^h), \forall \chi^h \in V_h, (a) \\ (\theta_x, w_x^h) = (\delta, w_x^h) + (\gamma, w_x^h), \forall w^h \in V_h, (b) \\ (\xi_t, \phi^h) + (\xi_x, \phi_x^h) + 2(uq - u^h q^h, \phi_x^h) \\ - (qv - q^h v^h, \phi_x^h) - (u\sigma - u^h \sigma^h, \phi_x^h) \\ = -(\rho_t, \phi^h) + \lambda(\rho, \phi^h), \forall \phi^h \in W_h, (c) \\ (\gamma_t, \psi^h) + (\gamma_x, \psi_x^h) + 2(v\sigma - v^h \sigma^h, \psi_x^h) \\ - (\sigma u - \sigma^h u^h, \psi_x^h) - (qv - q^h v^h, \psi_x^h) \\ = -(\delta_t, \psi^h) + \lambda(\delta, \psi^h), \forall \psi^h \in W_h, (d) \end{cases} \quad (11)$$

Theorem 3.1: Assuming that $u^h(0) = \tilde{u}^h(0), v^h(0) = \tilde{v}^h(0), q^h(0) = \tilde{q}^h(0), \sigma^h(0) = \tilde{\sigma}^h(0)$, we have

$$\|u - u^h\|^2 + h^2 \|u - u^h\|_1^2 \leq Ch^{2\min(k+1, r+1)},$$

$$\|v - v^h\|^2 + h^2 \|v - v^h\|_1^2 \leq Ch^{2\min(k+1, r+1)},$$

$$\|q - q^h\|^2 + h^2 \|q - q^h\|_1^2 \leq Ch^{2\min(k+1, r+1)},$$

$$\|\sigma - \sigma^h\|^2 + h^2 \|\sigma - \sigma^h\|_1^2 \leq Ch^{2\min(k+1, r+1)}.$$

Proof: Since estimates of ρ, δ, η and τ are given, respectively, it is sufficient to estimate ξ, γ, ς and θ . Choosing $\chi^h = \varsigma$ in (11a) and $w^h = \theta$ in (11b), using the Cauchy-Schwarz's inequality and Young's inequality, we have

$$\|\varsigma_x\| \leq C(\|\rho\| + \|\xi\|), \|\theta_x\| \leq C(\|\delta\| + \|\gamma\|). \quad (12)$$

Using Poincaré inequality, we have

$$\|\varsigma\| \leq C(\|\rho\| + \|\xi\|), \|\theta\| \leq C(\|\delta\| + \|\gamma\|). \quad (13)$$

Take $\phi^h = \xi$ in (11c) and use (13) to obtain

$$\begin{aligned} & (\xi_t, \xi) + A(\xi, \xi) \\ & \leq C(\|\rho\|^2 + \|\rho_t\|^2 + \|\xi\|^2) \\ & \quad + C\|q\|_{0,\infty}(\|\eta\| + \|\rho\| + \|\xi\| + \|\tau\| + \|\delta\| + \|\gamma\|)\|\xi\|_1 \\ & \quad + C\|u^h\|_{0,\infty}(\|\rho\| + \|\xi\| + \|\delta\| + \|\gamma\|)\|\xi\|_1 \\ & \quad + C\|\sigma\|_{0,\infty}(\|\eta\| + \|\rho\| + \|\xi\|)\|\xi\|_1 \\ & \quad + C\|v^h\|_{0,\infty}(\|\rho\| + \|\xi\|)\|\xi\|_1. \end{aligned} \quad (14)$$

Take $\psi^h = \gamma$ in (11d) and use (13) to have

$$\begin{aligned} & (\gamma_t, \gamma) + A(\gamma, \gamma) \\ & \leq C(\|\delta\|^2 + \|\delta_t\|^2 + \|\gamma\|^2) \\ & \quad + C\|\sigma\|_{0,\infty}(\|\tau\| + \|\delta\| + \|\gamma\| + \|\eta\| + \|\rho\| + \|\xi\|)\|\gamma\|_1 \\ & \quad + C\|v^h\|_{0,\infty}(\|\delta\| + \|\gamma\| + \|\rho\| + \|\xi\|)\|\gamma\|_1 \\ & \quad + C\|q\|_{0,\infty}(\|\tau\| + \|\delta\| + \|\gamma\|)\|\gamma\|_1 \\ & \quad + C\|u^h\|_{0,\infty}(\|\tau\| + \|\delta\| + \|\gamma\|)\|\gamma\|_1. \end{aligned} \quad (15)$$

Add (14) and (15) to get

$$\begin{aligned} & \frac{d}{dt} [\|\xi\|^2 + \|\gamma\|^2] + A(\xi, \xi) + A(\gamma, \gamma) \\ & \leq C(\|\xi\|^2 + \|\gamma\|^2) + C(\|\rho\|^2 + \|\rho_t\|^2 + \|\delta\|^2 \\ & \quad + \|\delta_t\|^2 + \|\eta\|^2 + \|\tau\|^2) + \frac{\mu_0}{2} (\|\xi\|_1^2 + \|\gamma\|_1^2). \end{aligned} \quad (16)$$

Integrate (16) with respect to time t to get

$$\begin{aligned} & \|\xi\|^2 + \|\gamma\|^2 + \int_0^t [A(\xi, \xi) + A(\gamma, \gamma)] ds \\ & \leq C \int_0^t (\|\xi\|^2 + \|\gamma\|^2) ds + C \int_0^t (\|\rho\|^2 + \|\rho_t\|^2 + \|\delta\|^2 \\ & \quad + \|\delta_t\|^2 + \|\eta\|^2 + \|\tau\|^2) ds + \frac{\mu_0}{2} \int_0^t (\|\xi\|_1^2 + \|\gamma\|_1^2) ds. \end{aligned} \quad (17)$$

For (17), we use the Gronwall lemma and $A(w, w) \geq \mu_0 \|\xi\|_1^2$ to get

$$\begin{aligned} & \|\xi\|^2 + \|\gamma\|^2 + \int_0^t (\|\xi\|_1^2 + \|\gamma\|_1^2) ds \\ & \leq C \int_0^t (\|\rho\|^2 + \|\rho_t\|^2 + \|\delta\|^2 + \|\delta_t\|^2 + \|\eta\|^2 + \|\tau\|^2) ds. \end{aligned} \quad (18)$$

Choosing $\phi^h = \xi_t$ in (11c) and using (13), we have

$$\begin{aligned} & (\xi_t, \xi_t) + \frac{1}{2} \frac{d}{dt} A(\xi, \xi) \\ & \leq C(\|\rho_t\|^2 + \|\rho\|^2 + \|\xi\|^2) + C\|u^h\|_{0,\infty}(\|\rho\| + \|\xi\| \\ & \quad + \|\delta\| + \|\gamma\|)\|\xi_t\|_1 + C\|q\|_{0,\infty}(\|\eta\| + \|\rho\| + \|\xi\| + \\ & \quad \|\tau\| + \|\delta\| + \|\gamma\|)\|\xi_t\|_1 + C\|v^h\|_{0,\infty}(\|\rho\| + \|\xi\|)\|\xi_t\|_1 \\ & \quad + C\|\sigma\|_{0,\infty}(\|\eta\| + \|\rho\| + \|\xi\|)\|\xi_t\|_1 + \mu_1 \|\xi_t\|^2 \end{aligned} \quad (19)$$

Take $\psi^h = \gamma_t$ in (11d) to get

$$\begin{aligned} & (\gamma_t, \gamma_t) + \frac{1}{2} \frac{d}{dt} A(\gamma, \gamma) \\ & \leq C(\|\delta_t\|^2 + \|\delta\|^2 + \|\gamma\|^2) + C\|\sigma\|_{0,\infty}(\|\eta\| + \|\varsigma\| \\ & \quad + \|\tau\| + \|\theta\|)\|\gamma_t\|_1 + C\|v^h\|_{0,\infty}(\|\rho\| + \|\xi\| \\ & \quad + \|\delta\| + \|\gamma\|)\|\gamma_t\|_1 + C\|u^h\|_{0,\infty}(\|\delta\| + \|\gamma\|)\|\gamma_t\|_1 \\ & \quad + C\|q\|_{0,\infty}(\|\tau\| + \|\theta\|)\|\gamma_t\|_1 + \mu_2\|\gamma_t\|^2 \\ & \leq C(\|\delta_t\|^2 + \|\delta\|^2 + \|\gamma\|^2) + C\|\sigma\|_{0,\infty}(\|\eta\| + \|\rho\| \\ & \quad + \|\xi\| + \|\tau\| + \|\delta\| + \|\gamma\|)\|\gamma_t\|_1 + C\|v^h\|_{0,\infty}(\|\rho\| + \|\xi\| \\ & \quad + \|\delta\| + \|\gamma\|)\|\gamma_t\|_1 + C\|u^h\|_{0,\infty}(\|\delta\| + \|\gamma\|)\|\gamma_t\|_1 \\ & \quad + C\|q\|_{0,\infty}(\|\tau\| + \|\delta\| + \|\gamma\|)\|\gamma_t\|_1 + \mu_2\|\gamma_t\|^2. \end{aligned} \quad (20)$$

Add (19) and (20), integrate with respect to time t and use the Gronwall lemma to get

$$\begin{aligned} & \mu_3 \int_0^t [\|\xi\|^2 + \|\gamma\|^2] ds + \|\xi\|_1^2 + \|\gamma\|_1^2 \\ & \leq C \int_0^t (\|\rho\|^2 + \|\rho_t\|^2 + \|\delta\|^2 + \|\delta_t\|^2 + \|\eta\|^2 + \|\tau\|^2) ds. \end{aligned} \quad (21)$$

Combining (12), (13), (18) and (21), we have

$$\begin{aligned} \|\varsigma\|^2 & \leq \|\varsigma\|_1^2 \leq C(\|\rho\|^2 + \|\xi\|^2) \\ & \leq C\|\rho\|^2 + C \int_0^t (\|\rho\|^2 + \|\rho_t\|^2 + \|\delta\|^2 \\ & \quad + \|\delta_t\|^2 + \|\eta\|^2 + \|\tau\|^2) ds. \end{aligned} \quad (22)$$

$$\begin{aligned} \|\theta\|^2 & \leq \|\theta\|_1^2 \leq C(\|\delta\|^2 + \|\gamma\|^2) \\ & \leq C\|\delta\|^2 + C \int_0^t (\|\rho\|^2 + \|\rho_t\|^2 + \|\delta\|^2 \\ & \quad + \|\delta_t\|^2 + \|\eta\|^2 + \|\tau\|^2) ds. \end{aligned} \quad (23)$$

Using (18), (21)-(23), (7)-(10) with the triangle inequality, we obtain the conclusion. ■

IV. FULLY-DISCRETE ERROR ESTIMATES

For the backward Euler procedure, let $0 = t_0 < t_1 < \dots < t_M = T$ be a given partition of the time interval $[0, T]$ with step length $\Delta t = T/M$, for some positive integer M . For a smooth function ϕ on $[0, T]$, define $\phi^n = \phi(t_n)$ and $\bar{\partial}_t \phi^n = (\phi^n - \phi^{n-1})/\Delta t$.

Let U^n , V^n , Q^n and Z^n , respectively, be the approximations of u , v , q and σ at $t = t_n$ which we shall define through the following scheme. Given $\{U^{n-1}, V^{n-1}, Q^{n-1}, Z^{n-1}\}$ in $V_h \times W_h$, we now determine $\{U^n, V^n, Q^n, Z^n\}$ in $V_h \times W_h$ satisfying

$$\begin{cases} (U_x^n, \chi_x^h) = (Q^h, \chi_x^h), \forall \chi^h \in V_h, (a) \\ (V_x^n, w_x^h) = (Z^n, w_x^h), \forall w^h \in V_h, (b) \\ (\bar{\partial}_t Q^n, \phi^h) + (Q_x^n, \phi_x^h) + 2(U^n Q^n, \phi_x^h) - (Q^n V^n, \phi_x^h) \\ - (U^n Z^n, \phi_x^h) = -(f^n, \phi_x^h), \forall \phi^h \in W_h, (c) \\ (\bar{\partial}_t Z^n, \psi^h) + (Z_x^n, \psi_x^h) + 2(V^n Z^n, \psi_x^h) - (Z^n U^n, \psi_x^h) \\ - (Q^n V^n, \psi_x^h) = -(g^n, \psi_x^h), \forall \psi^h \in W_h, (d) \end{cases} \quad (24)$$

we now split the errors

$$\begin{aligned} u(t_n) - U^n &= u(t_n) - \tilde{u}^h(t_n) + \tilde{u}^h(t_n) - U^n = \eta^n + \varsigma^n \\ v(t_n) - V^n &= v(t_n) - \tilde{v}^h(t_n) + \tilde{v}^h(t_n) - V^n = \tau^n + \theta^n \\ q(t_n) - Q^n &= q(t_n) - \tilde{q}^h(t_n) + \tilde{q}^h(t_n) - Q^n = \rho^n + \xi^n \\ \sigma(t_n) - Z^n &= \sigma(t_n) - \tilde{\sigma}^h(t_n) + \tilde{\sigma}^h(t_n) - Z^n = \sigma^n + \gamma^n \end{aligned}$$

Using (5)-(6) and (24), we obtain the following error equation

$$\begin{cases} (\varsigma_x^n, \chi_x^h) = (\rho^n + \xi^n, \chi_x^h), \forall \chi^h \in V_h, (a) \\ (\theta_x^n, w_x^h) = (\delta^n + \gamma^n, w_x^h), \forall w^h \in V_h, (b) \\ (\bar{\partial}_t \xi^n, \phi^h) + (\xi_x^n, \phi_x^h) + 2(u^n q^n - U^n Q^n, \phi_x^h) \\ - (q^n v^n - Q^n V^n, \phi_x^h) - (u^n \sigma^n - U^n Z^n, \phi_x^h) \\ = -(\pi^n + \bar{\partial}_t \rho^n, \phi^h) + \lambda(\rho^n, \phi^h), \forall \phi^h \in W_h, (c) \\ (\bar{\partial}_t \gamma^n, \psi^h) + (\gamma_x^n, \psi_x^h) + 2(v^n \sigma^n - V^n Z^n, \psi_x^h) \\ - (\sigma^n u^n - Z^n U^n, \psi_x^h) - (q^n v^n - Q^n V^n, \psi_x^h) \\ = -(\varepsilon^n + \bar{\partial}_t \delta^n, \psi^h) + \lambda(\delta^n, \psi^h), \forall \psi^h \in W_h, (d) \end{cases} \quad (25)$$

where $\pi^n = q_t(t_n) - \bar{\partial}_t q(t_n)$, $\varepsilon^n = \sigma_t(t_n) - \bar{\partial}_t \sigma(t_n)$.

Theorem 4.1: With $Q^0(0) = \bar{q}^h(0)$, $Z^0(0) = \bar{\sigma}^h(0)$ and $1 \leq J \leq M$, we have

$$\begin{aligned} & \|q^J - Q^J\| + \|\sigma^J - Z^J\| + \|u^J - U^J\| + \|v^J - V^J\| \\ & + h\|u^J - U^J\|_1 + h\|v^J - V^J\|_1 \leq C h^{min(r+1, k+1)} \end{aligned}$$

Proof: Take $\chi^h = \varsigma^n$ and $w^h = \theta^n$ in (25a,b) to get

$$\|\varsigma_x^n\| \leq C(\|\rho^n\| + \|\xi^n\|), \|\theta_x^n\| \leq C(\|\delta^n\| + \|\gamma^n\|). \quad (26)$$

Using the Poincare inequality, we have

$$\|\varsigma^n\| \leq C(\|\rho^n\| + \|\xi^n\|), \|\theta^n\| \leq C(\|\delta^n\| + \|\gamma^n\|). \quad (27)$$

Choose $\phi^h = \xi^n$ and $\psi^h = \gamma^n$ in (25c,d) and add the two equations to get

$$\begin{aligned} & (\bar{\partial}_t \xi^n, \xi^n) + (\bar{\partial}_t \gamma^n, \gamma^n) + (\xi_x^n, \xi_x^n) + (\gamma_x^n, \gamma_x^n) \\ & = -2(u^n q^n - U^n Q^n, \xi_x^n) + (q^n v^n - Q^n V^n, \xi_x^n) \\ & \quad + (u^n \sigma^n - U^n Z^n, \xi_x^n) - (\pi^n + \bar{\partial}_t \rho^n, \xi^n) + \lambda(\rho^n, \xi^n) \\ & \quad - 2(v^n \sigma^n - V^n Z^n, \gamma_x^n) + (\sigma^n u^n - Z^n U^n, \gamma_x^n) \\ & \quad + (q^n v^n - Q^n V^n, \gamma_x^n) - (\varepsilon^n + \bar{\partial}_t \delta^n, \gamma^n) + \lambda(\delta^n, \gamma^n). \end{aligned} \quad (28)$$

Noting that $(\bar{\partial}_t \xi^n, \xi^n) \geq \frac{1}{2} \bar{\partial}_t \|\xi^n\|^2$, $(\bar{\partial}_t \gamma^n, \gamma^n) \geq \frac{1}{2} \bar{\partial}_t \|\gamma^n\|^2$, and using (27), we have

$$\begin{aligned} & \frac{1}{2} \bar{\partial}_t \|\xi^n\|^2 + \frac{1}{2} \bar{\partial}_t \|\gamma^n\|^2 + \|\xi_x^n\|^2 + \|\gamma_x^n\|^2 \\ & \leq C(\|\rho^n\|^2 + \|\eta^n\|^2 + \|\delta^n\|^2 + \|\tau^n\|^2) + C(\|\pi^n\|^2 \\ & \quad + \|\varepsilon^n\|^2) + C(\bar{\partial}_t \|\rho^n\|^2 + \bar{\partial}_t \|\delta^n\|^2) + C(\|\xi^n\|^2 + \|\gamma^n\|^2). \end{aligned} \quad (29)$$

Noting that $\bar{\partial}_t \|\xi^n\|^2 = (\|\xi^n\|^2 - \|\xi^{n-1}\|^2)/\Delta t$, $\bar{\partial}_t \|\gamma^n\|^2 = (\|\gamma^n\|^2 - \|\gamma^{n-1}\|^2)/\Delta t$, and combining (29), we have

$$\begin{aligned} & \|\xi^n\|^2 - \|\xi^{n-1}\|^2 + \|\gamma^n\|^2 - \|\gamma^{n-1}\|^2 \\ & + 2\Delta t(\|\xi_x^n\|^2 + \|\gamma_x^n\|^2) \leq C\Delta t(\|\rho^n\|^2 + \|\eta^n\|^2 \\ & + \|\delta^n\|^2 + \|\tau^n\|^2) + C\Delta t(\|\pi^n\|^2 + \|\varepsilon^n\|^2) \\ & + C\Delta t(\bar{\partial}_t \|\rho^n\|^2 + \bar{\partial}_t \|\delta^n\|^2) + C\Delta t(\|\xi^n\|^2 + \|\gamma^n\|^2). \end{aligned} \quad (30)$$

Sum (30) from $n = 1$ to J ($1 \leq J \leq M$) and use Gronwall lemma to get

$$\begin{aligned} & (1 - C\Delta t)(\|\xi^J\|^2 + \|\gamma^J\|^2) + 2\Delta t \sum_{n=1}^J (\|\xi_x^n\|^2 + \|\gamma_x^n\|^2) \\ & \leq C(\|\xi^0\|^2 + \|\gamma^0\|^2) + C\Delta t \sum_{n=1}^J (\|\rho^n\|^2 + \|\eta^n\|^2 \\ & \quad + \|\delta^n\|^2 + \|\tau^n\|^2) + C\Delta t \sum_{n=1}^J (\|\pi^n\|^2 + \|\varepsilon^n\|^2) \\ & \quad + C\Delta t \sum_{n=1}^J (\bar{\partial}_t \|\rho^n\|^2 + \bar{\partial}_t \|\delta^n\|^2). \end{aligned} \quad (31)$$

Use

$$\bar{\partial}_t \|\rho^n\|^2 \leq \frac{h^{2(r+1)}}{\Delta t} \int_{t_{n-1}}^{t_n} \|q_t\|_{r+1}^2 ds,$$

$$\bar{\partial}_t \|\delta^n\|^2 \leq \frac{h^{2(r+1)}}{\Delta t} \int_{t_{n-1}}^{t_n} \|\sigma_t\|_{r+1}^2 ds$$

$$\|\pi^n\|^2 \leq C\Delta t \int_{t_{n-1}}^{t_n} \|q_{tt}\|_{r+1}^2 ds,$$

$$\|\varepsilon^n\|^2 \leq C\Delta t \int_{t_{n-1}}^{t_n} \|\sigma_{tt}\|_{r+1}^2 ds,$$

and (7)-(10) to obtain

$$\begin{aligned} & (\|\xi^J\|^2 + \|\gamma^J\|^2) + 2\Delta t \sum_{n=1}^J (\|\xi_x^n\|^2 + \|\gamma_x^n\|^2) \\ & \leq Ch^{2\min(r+1, k+1)} (\|u\|_{L^\infty(H^{k+1})}^2 + \|u_t\|_{L^\infty(H^{k+1})}^2 \\ & \quad + \|v\|_{L^\infty(H^{k+1})}^2 + \|v_t\|_{L^\infty(H^{k+1})}^2 + \|q\|_{L^\infty(H^{r+1})}^2 \\ & \quad + \|\sigma\|_{L^\infty(H^{r+1})}^2 + \|q_t\|_{L^2(H^{k+1})}^2 + \|\sigma_t\|_{L^2(H^{k+1})}^2) \\ & \quad + C\Delta t^2 (\|q_{tt}\|_{L^2(L^2)}^2 + \|\sigma_{tt}\|_{L^2(L^2)}^2). \end{aligned} \quad (32)$$

Using (26)-(27), we have

$$\begin{aligned} & \|\varsigma^J\|_1 + \|\theta^J\|_1 \\ & \leq Ch^{r+1} (\|q\|_{L^\infty(H^{r+1})} + \|\sigma\|_{L^\infty(H^{r+1})}) \\ & \quad + Ch^{\min(r+1, k+1)} (\|u\|_{L^\infty(H^{k+1})} + \|u_t\|_{L^\infty(H^{k+1})} \\ & \quad + \|v\|_{L^\infty(H^{k+1})} + \|v_t\|_{L^\infty(H^{k+1})} + \|q\|_{L^\infty(H^{r+1})} \\ & \quad + \|\sigma\|_{L^\infty(H^{r+1})} + \|q_t\|_{L^2(H^{k+1})} + \|\sigma_t\|_{L^2(H^{k+1})}) \\ & \quad + C\Delta t (\|q_{tt}\|_{L^2(L^2)} + \|\sigma_{tt}\|_{L^2(L^2)}). \end{aligned} \quad (33)$$

We apply the triangle inequality to get the conclusion. ■

REFERENCES

- [1] R C Mittal, G Arora. Numerical solution of the coupled viscous Burgers' equation, Communications in Nonlinear Science and Numerical Simulation, 2010, 16(3): 1304-1313.
- [2] A K Pani. An H^1 -Galerkin mixed finite element methods for parabolic partial differential equations, SIAM J. Numer. Anal., 1998, 35: 712-727.
- [3] A K Pani, G Fairweather. H^1 -Galerkin mixed finite element methods for parabolic partial integro-differential equations, IMA Journal of Numerical Analysis., 2002, 22: 231-252.
- [4] D Y Shi, H H Wang. An H^1 -Galerkin nonconforming mixed finite element method for integro-differential equation of parabolic type, Journal of Mathematical Research & Exposition, 2009, 29(5): 871-881.
- [5] R W Wang. Error estimates for H^1 -Galerkin mixed finite element methods hyperbolic type integro- differential equation, Math. Numer. Sin., 2006, 28(1): 20-30. (in Chinese)
- [6] Y Liu, H Li, S He. Error estimates of H^1 -Galerkin mixed finite element methods for pseudo-hyperbolic partial integro-differential equation, Numerical Mathematics A Journal of Chinese Universities, 2010, 32(1): 1-20. (in Chinese)
- [7] L Guo, H Z Chen. H^1 -Galerkin mixed finite element methods for the Sobolev equation, Journal of Systems Science and Mathematical Sciences, 2006, 26(3): 301-314. (in Chinese)
- [8] D Y Shi, H H Wang. Nonconforming H^1 -Galerkin mixed FEM for Sobolev equations on anisotropic Meshes, Acta Mathematicae Applicatae Sinica (English Series), 2009, 25(2): 335-344.
- [9] Y Liu, H Li. H^1 -Galerkin mixed finite element methods for pseudo-hyperbolic equations, Appl. Math. Comput., 2009, 212: 446-457.
- [10] Y Liu, J F Wang, H Li, W Gao, S He. A new splitting H^1 -Galerkin mixed method for pseudo-hyperbolic equations, International Journal of Engineering and Natural Sciences, 2011, 5(2): 58-63.
- [11] Z J Zhou. An H^1 -Galerkin mixed finite element method for a class of heat transport equations, Applied Mathematical Modelling, 2010, 34(9): 2414-2425.
- [12] J F Wang, Y Liu, H Li, X Y Li. H^1 -Galerkin mixed element method for the coupling nonlinear parabolic partial equations, Pure Mathematics, 2011, 1(2): 73-79. (in Chinese)
- [13] Y Liu, H Li, J F Wang. Error estimates of H^1 -Galerkin mixed finite element method for Schrödinger equation. Appl. Math. J. Chinese Univ. 2009, 24(1): 83-89.
- [14] Y Liu, H Li. A new mixed finite element method for pseudo-hyperbolic equation, Mathematica Applicata, 2010, 23(1): 150-157.
- [15] Y. Liu. Analysis and numerical simulation of nonstandard mixed element methods, PhD thesis, Inner Mongolia University, Hohhot, China, 2011.
- [16] A K Pani, R K Sinha, A K Otta. An H^1 -Galerkin mixed method for second order hyperbolic equations, Inter Journal of Numerical Anal and Modeling., 2004, 1(2): 111-129.
- [17] H Z Chen, H Wang. An optimal-order error estimate on an H^1 -Galerkin mixed method for a nonlinear parabolic equation in porous medium flow, Numer. Methods Partial Differential Equations, 2010, 26: 188-205.
- [18] H T Che, Y J Wang, Z J Zhou. An optimal error estimates of H^1 -Galerkin expanded mixed finite element methods for nonlinear viscoelasticity-type equation, Mathematical Problems in Engineering, Volume 2011, Article ID 570980, 18 pages. doi:10.1155/2011/570980.
- [19] M F Wheeler. A priori L^2 -error estimates for Galerkin approximations to parabolic differential equation, SIAM J. Numer. Anal., 1973, 10: 723-749.

ACKNOWLEDGMENT

This work is supported by National Natural Science Fund (No. 11061021), the Scientific Research Projection of Higher Schools of Inner Mongolia (No. NJ10006) and YSF of Inner Mongolia University (No. ND0702)