# Probing Anomalous $W W \gamma$ and $W W Z$ Couplings with Polarized Electron Beam at the LHeC and FCC-Ep Collider 

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#### Abstract

We study the anomalous $W W \gamma$ and $W W Z$ couplings by calculating total cross sections of two processes at the LHeC with electron beam energy $\mathrm{Ee}=140 \mathrm{GeV}$ and the proton beam energy $\mathrm{Ep}=7$ TeV , and at the FCC-ep collider with the polarized electron beam energy $\mathrm{Ee}=80 \mathrm{GeV}$ and the proton beam energy $\mathrm{Ep}=50 \mathrm{TeV}$. At the LHeC with electron beam polarization, we obtain the results for the difference of upper and lower bounds as $(0.975,0.118)$ and $(0.285$, 0.009 ) for the anomalous ( $\Delta \kappa \gamma, \lambda \gamma$ ) and ( $\Delta \kappa z, \lambda z$ ) couplings, respectively. As for FCC-ep collider, these bounds are obtained as (1.101, 0.065) and $(0.320,0.002)$ at an integrated luminosity of $\mathrm{L}_{\mathrm{int}}=100 \mathrm{fb}^{-1}$.


Keywords-Anomalous Couplings, Future Circular Collider, Large Hadron electron Collider, $W$-boson and $Z$-boson.

## I. Introduction

THE $S U(2) \times U(1)$ gauge symmetry of the Standard Model (SM) results in the triple gauge boson interactions. A precise determination of the trilinear gauge boson couplings is necessary to test the validity of the SM and the presence of new physics up to a high energy scale. Since the tree-level couplings of the $W W \gamma$ and $W W Z$ vertices are fixed by the SM, any deviations from their SM values would indicate the new physics beyond the SM. The photoproduction of the $W$ and $Z$ bosons through triple gauge boson interactions in the leptonhadron colliders HERA+LC and in the Large Hadron electron Collider ( LHeC ) has been studied theoretically in the papers [1]-[3] and [4], respectively. An investigation of the potential of the LHeC to probe anomalous $W W \gamma$ coupling has been presented in [5], [6].

The present bounds on the anomalous $W W \gamma$ and $W W Z$ couplings are provided by the LEP [7], Tevatron [8], [9] and LHC [10], [11] experiments.

Recently, the ATLAS [10], [11] and CMS [12], [13] Collaborations have established updated constraints on the anomalous $\mathrm{WW} \gamma$ and $W W Z$ couplings from the $\gamma W(Z)$ and $W^{+} W$ production processes. The results from ATLAS and
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CMS experiments based on two-parameter analysis of the anomalous couplings are given in Table I.
In this work, we investigate the $e p \rightarrow v e^{q \gamma X}$ and $e p \rightarrow v e^{q Z X}$ processes with anomalous $W W \gamma$ and $W W Z$ couplings at the high energy electron-proton collider LHeC and FCC-ep (Future Circular Collider-electron proton) collider [14]. LHeC is considered to be realised by accelerating electrons 140 GeV and colliding them with the 7 TeV protons. We take into account the energies of the FCC-ep as 80 GeV for electron beam and 50 TeV for proton beam. We also consider the possibility of the electron beam polarization at LHeC [15] and FCC-ep which extends the sensitivity to anomalous triple gauge boson couplings.

TABLE I
The Available 95\% C.L. Two-Parameter Bounds on Anomalous
Couplings ( $\Delta \kappa \gamma, \lambda \gamma$ ) and ( $\Delta K z, \lambda Z$ ) FROM the Atlas and CMS EXPERIMENTS

|  | ATLAS | CMS | ATLAS (upper- <br> lower) | CMS (upper- <br> lower) |
| :---: | :---: | :---: | :---: | :---: |
| $\Delta \kappa \gamma$ | $-0.420,0.480$ | $-0.250,0.250$ | 0.900 | 0.500 |
| $\lambda \gamma$ | $-0.068,0.062$ | $-0.050,0.042$ | 0.130 | 0.092 |
| $\Delta \kappa z$ | $-0.045,0.045$ | $-0.160,0.180$ | 0.090 | 0.340 |
| $\lambda \mathrm{z}$ | $-0.063,0.063$ | $-0.055,0.055$ | 0.126 | 0.110 |

## II. ANOMALOUS COUPLINGS

The $W W \gamma$ and $W W Z$ interaction vertices are described by an effective Lagrangian with the coupling constants $\mathrm{g}_{\mathrm{wW}}^{\gamma}$ and $g_{W W Z}$ and dimensionless parameter pairs ( $\left.\Delta \kappa \gamma, \lambda \gamma\right)$ and ( $\Delta \kappa z$, $\lambda z)$

$$
\begin{align*}
& L=i g w w_{\gamma}\left[g_{1}^{\gamma}\left(W_{\mu \nu}^{\dagger} W^{\mu} A^{v}-W^{\mu \nu} W_{\mu}^{\dagger} A_{\nu}\right)+\kappa_{\gamma} W_{\mu}^{\dagger} W_{v} A^{\mu v}+\frac{\lambda_{v}}{m_{W}^{2}} W_{\rho \mu}^{\dagger} W_{v}^{\mu} A^{v \rho}\right]+ \\
& i g w w_{Z}\left[g_{1}^{Z}\left(W_{\mu \nu}^{\dagger} W^{\mu} Z^{v}-W^{\mu \nu} W_{\mu}^{\dagger} Z_{v}\right)+\kappa_{Z} W_{\mu}^{\dagger} W_{v} Z^{\mu v}+\frac{\lambda_{Z}}{m_{W}^{2}} W_{\rho \mu}^{\dagger} W_{v}^{\mu} Z^{v \rho}\right] \tag{1}
\end{align*}
$$

where $g_{w W \gamma}=g_{e}=g \sin \theta_{W}$ and $g_{w W Z}=g \cos \theta_{W}$. In general these vertices involve six C and P conserving couplings [16]. However, the electromagnetic gauge invariance requires that $g_{1}^{\gamma}=1$. The anomalous couplings are defined as $\kappa_{V}=1+\Delta \kappa_{V}$ where $V=\gamma, Z$ and $g_{1}^{Z}=1+\Delta g_{1}^{Z}$. The $W_{\mu \nu}, Z_{\mu \nu}$ and $A_{\mu \nu}$ are the field strength tensors for the $W$ - boson, $Z$ - boson and photon, respectively.
The one-loop corrections to the $W W \gamma$ and $W W Z$ vertices within the framework of the SM have been studied in [17][19]. These corrections to the $\Delta \kappa_{V}$ and $\lambda_{V}$ have been found to be of the order of $10^{-2}$ and $10^{-3}$, respectively. The values of the
couplings $\kappa_{\gamma}=\kappa_{Z}=1$ and $\lambda_{\gamma}=\lambda_{Z}=0$ correspond to the case of the SM. Since unitarity restricts the $W W \gamma$ and $W W Z$ couplings to their SM values at very high energies, the triple gauge couplings are modified as $\Delta \kappa_{V}\left(q^{2}\right)=\Delta \kappa_{V}(0) /\left(1+q^{2} / \Lambda^{2}\right)^{2}$ and $\lambda_{V}\left(q^{2}\right)=\lambda_{V}(0) /\left(1+q^{2} / \Lambda^{2}\right)^{2}$ where $V=\gamma, Z$. The $q^{2}$ is the square of momentum transfer into the process and $\Lambda$ is the new physics energy scale. The $\Delta \kappa_{V}(0)$ and $\lambda_{V}(0)$ are the values of the anomalous couplings at $q^{2}=0$. We assume the values of the anomalous couplings remain approximate constant in the interested energy scale $\left.\left(\Lambda^{2}>q\right)^{2}\right)$. We take $\Delta \kappa_{V}$ and $\lambda_{V}$ as free parameters in the considered range and find the bounds on these couplings effectively. For the numerical calculations, we have implemented interactions terms in the CalcHEP [20].


Fig. 1 Representative Feynman diagrams for subprocess $e q \rightarrow v e^{\gamma q^{\prime}}$


Fig. 2 Representative Feynman diagrams for subprocess $e q \rightarrow v e^{Z q^{\prime}}$

## III. Production Cross Sections for LHEC

According to the effective Lagrangian, the anomalous vertices for triple gauge interactions $W W \gamma$ and $W W Z$ are presented in the Feynman graphs as shown in Figs. 1 and 2. In order to calculate the cross sections for the process $e p \rightarrow v{ }_{e} q \gamma X$ and $e p \rightarrow v e^{q Z X}$, we apply the transverse momentum cut on photon and jet as $p_{T}^{\gamma}>50 \mathrm{GeV}, p_{T}^{j}>20 \mathrm{GeV}$; missing transverse momentum cut $p_{T}^{v}>20 \mathrm{GeV}$, pseudorapidity cuts $\left|\eta_{\gamma, j}\right|<3.5$; a cone radius cut between photons and jets $\Delta R_{\gamma, j}>$ 1.5. Using these cuts and the parton distribution functions of CTEQ6L [21], the total cross sections of the process $e p \rightarrow v e^{q \gamma X}$ as a function of anomalous couplings $\Delta \kappa_{\gamma}$ and $\lambda_{\gamma}$ for $E_{e}=140$ GeV with electron beam polarizations $P_{e}= \pm 0.8$ and $P_{e}=0$ are
presented in Figs. 3 and 4. In Figs. 5 and 6, the total cross sections of the $e p \rightarrow v e^{q Z X}$ process are given for the same energy. It is clear from these figures that the polarization ( $P_{e}=-0.8$ ) enhances the cross sections according to the unpolarized case.


Fig. 3 The cross section depending on anomalous coupling $\Delta \kappa \gamma$ of the process $e p \rightarrow v e^{q \gamma X}$ at $\mathrm{Ee}=140 \mathrm{GeV}$ for different electron beam b polarizations


Fig. 4 The cross section depending on anomalous coupling $\lambda \gamma$ of the process $e p \rightarrow v e^{q \gamma X}$ at $\mathrm{Ee}=140 \mathrm{GeV}$ for different electron beam polarizations


Fig. 5 The cross section depending on anomalous $\Delta \kappa_{Z}$ coupling of the process $e p \rightarrow v_{e} q Z X$ for $E_{e}=140 \mathrm{GeV}$


Fig. 6 The cross section depending on anomalous $\lambda_{Z}$ coupling of the process $e p \rightarrow v e^{q Z X}$ for $E_{e}=140 \mathrm{GeV}$

## IV. ANALYSIS FOR LHEC

In order to estimate the sensitivity to the anomalous $W W \gamma$ and $W W Z$ couplings, we use the $\chi^{2}$ function:

$$
\begin{equation*}
\chi^{2}\left(\Delta \kappa_{V}, \lambda_{V}\right)=\left(\frac{\sigma_{S M}-\sigma\left(\Delta \kappa_{V}, \lambda_{V}\right)}{\Delta \sigma_{S M}}\right)^{2} \tag{2}
\end{equation*}
$$

where $\quad \Delta \sigma_{S M}=\sigma_{S M} \sqrt{\delta_{\text {stat. }}^{2}} \quad$ with $\quad \delta_{\text {stat. }}=1 / \sqrt{N_{S M}} \quad$ and $N_{S M}=\sigma_{S M} L$. In our calculations, we consider that two of the couplings ( $\Delta \kappa, \lambda$ ) are assumed to deviate from their SM value. We estimate the sensitivity to the anomalous couplings at 95 C.L. at the LHeC for the integrated luminosities of $10 \mathrm{fb}^{-1}$ and $100 \mathrm{fb}^{-1}$. The contour plots of anomalous couplings in $\Delta \kappa_{\gamma}{ }^{-} \lambda_{\gamma}$ plane for the integrated luminosities of $10 \mathrm{fb}^{-1}$ and $100 \mathrm{fb}^{-1}$ at electron beam energies $E_{e}=140 \mathrm{GeV}$ are given in Fig. 7. The contour plots of anomalous couplings in $\Delta \kappa_{Z}{ }^{-\lambda}{ }_{Z}$ plane for the integrated luminosities of $10 \mathrm{fb}^{-1}$ and $100 \mathrm{fb}^{-1}$ at electron beam energies of $E_{e}=140 \mathrm{GeV}$ are shown in Fig. 8.


Fig. 7 Two dimensional 95\% C.L contour plot anomalous couplings in the $\lambda_{\gamma}-\Delta \kappa_{\gamma}$ plane for the integrated luminosity of $10 \mathrm{fb}^{-1}$ and 100 $\mathrm{fb}^{-1}$ at electron beam energy $E_{e}=140 \mathrm{GeV}$ with polarization $P_{e}=-0.8$


Fig. 8 Two-dimensional 95\% C.L contour plot of anomalous couplings in the $\lambda_{Z}-\Delta \kappa_{Z}$ plane for the integrated luminosity of $10 \mathrm{fb}^{-1}$ and $100 \mathrm{fb}^{-1}$ at electron beam energy $E_{e}=140 \mathrm{GeV}$ with polarization

$$
P e=-0.8
$$

The difference of the upper and lower bounds on the anomalous couplings $\Delta \kappa_{V}$ and $\lambda_{V}($ where $V=\gamma, Z)$ can be written as

$$
\begin{equation*}
\delta \Delta \kappa_{V}=\Delta \kappa_{V}^{\text {upper }}-\Delta \kappa_{V}^{\text {lower }}, \delta \lambda_{V}=\lambda_{V}^{\text {upper }}-\lambda_{V}^{\text {lower }} \tag{3}
\end{equation*}
$$

The current limits on anomalous couplings and the difference of the upper and lower bounds for electron beam energies of 140 GeV with integrated luminosities $L_{i n t}=10 \mathrm{fb}^{-1}$ and $100 \mathrm{fb}^{-1}$ at LHeC with the unpolarized (polarized) electron beam are given in Table II. We have obtained two-parameter limits on $\delta \Delta \kappa_{\gamma}$ and $\delta \lambda_{\gamma}$ which can be compared to the ATLAS and CMS results. However, the limits on $\delta \lambda_{Z}$ is found to be much more sensitive than the current limits.

TABLE II
The 95\% C.L. Current Limits on the Anomalous Couplings and the Difference of the Upper and Lower Bounds for Electron Beam ENERGY OF $\mathrm{E}_{\mathrm{E}}=140$ GEV WITH $\mathrm{L}_{\mathrm{INT}}=100 \mathrm{FB}^{-1}$ FOR POLARIZED AND

UnPOLARIZED ELECTRON BEAM

| $\mathbf{P e}$ | $\Delta \boldsymbol{\kappa} \boldsymbol{\gamma}$ | $\delta \Delta \boldsymbol{\kappa} \boldsymbol{\gamma}$ | $\boldsymbol{\lambda} \boldsymbol{\gamma}$ | $\delta \lambda \boldsymbol{\gamma}$ |
| :---: | :---: | :---: | :---: | :---: |
| -0.8 | $-0.182,0.793$ | 0.975 | $-0.039,0.079$ | 0.118 |
| 0 | $0.192,0.798$ | 0.990 | $-0.041,0.081$ | 0.122 |
| 0.8 | $0.251,0.844$ | 1.095 | $-0.047,0.086$ | 0.133 |
| $\mathbf{P e}$ | $\Delta \boldsymbol{\kappa} \mathbf{Z}$ | $\delta \Delta \boldsymbol{\kappa} \mathbf{Z}$ | $\boldsymbol{\lambda} \mathbf{z}$ | $\delta \boldsymbol{z} \mathbf{z}$ |
| -0.8 | $-0.143,0.142$ | 0.285 | $-0.001,0.008$ | 0.009 |
| 0 | $0.273,0.089$ | 0.362 | $-0.003,0.009$ | 0.012 |
| 0.8 | $0.253,0.215$ | 0.468 | $-0.004,0.010$ | 0.014 |

## V. Production Cross Sections for FCC-EP

For calculate the cross sections for the process $e p \rightarrow v{ }_{e} q \gamma X$ and $e p \rightarrow v e^{q Z X}$, we apply the transverse momentum cut on photon and jet as $p_{T}^{\gamma}>20 \mathrm{GeV}, p_{T}^{j}>20 \mathrm{GeV}$; missing transverse momentum cut $p_{T}^{v}>20 \mathrm{GeV}$, pseudorapidity cuts $\eta_{\gamma, j}$ the range of between -5 and 0 ; Using these cuts and the parton distribution functions of CTEQ6M [14], the total cross sections of the process $e p \rightarrow v \gamma q X$ as a function of anomalous couplings $\Delta \kappa_{\gamma}$ and $\lambda_{\gamma}$ for $E_{e}=80 \mathrm{GeV}$ with $\left(P_{e}= \pm 0.8\right)$ and
without ( $P_{e}=0$ ) electron beam polarization are presented in Figs. 9 and 10. It is clear from these figures that the polarization $\left(P_{e}=-0.8\right)$ enhances the cross sections according to the unpolarized case.


Fig. 9 The cross section depending on anomalous coupling $\Delta \kappa_{\gamma}$ of the process $e p \rightarrow v{ }_{e}{ }^{q \gamma X}$ at $E_{e}=80 \mathrm{GeV}$ for different electron beam polarizations


Fig. 10 The cross section depending on anomalous $\lambda_{\gamma}$ coupling of the process $e p \rightarrow v{ }_{e} q \gamma X$ for $E_{e}=80 \mathrm{GeV}$


Fig. 11 The cross section depending on anomalous $\Delta \kappa_{Z}$ coupling of the process $e p \rightarrow v e^{q Z X}$ for $E_{e}=80 \mathrm{GeV}$

The cross sections depending on anomalous couplings $\Delta \kappa_{Z}$ and $\lambda_{Z}$ of the process $e p \rightarrow v e^{q Z X}$ for $E_{e}=80 \mathrm{GeV}$ with $P_{e}= \pm 0.8$
and without $\left(P_{e}=0\right)$ electron beam polarization are presented in Figs. 11 and 12.


Fig. 12 The cross section depending on anomalous $\lambda_{Z}$ coupling of the

$$
\text { process } e p \rightarrow v e_{e} q Z X \text { for } E_{e}=80 \mathrm{GeV}
$$

## VI. Analysis for FCC-EP

The contour plots of anomalous couplings in $\Delta \kappa_{\gamma}-\lambda_{\gamma}$ plane for the integrated luminosities of $10 \mathrm{fb}^{-1}$ and $100 \mathrm{fb}^{-1}$ at electron beam energies $E_{e}=80 \mathrm{GeV}$ are given in Fig. 13. For the process $e p \rightarrow v{ }_{e} q Z X$, we make analysis of the signal and backgrounds when $Z$ decays leptonically, $Z \rightarrow l^{+} l^{-}$where $l=e, \mu$. The contour plots of anomalous couplings in $\Delta \kappa_{Z}-\lambda Z$ plane for the integrated luminosities of $10 \mathrm{fb}^{-1}$ and $100 \mathrm{fb}^{-1}$ at electron beam energies of $E_{e}=80 \mathrm{GeV}$ are presented in Fig. 14.
The difference of the upper and lower bounds on the anomalous couplings $\Delta \kappa_{V}$ and $\lambda_{V}$ (where $V=\gamma, Z$ ) can be written as

$$
\begin{equation*}
\delta \Delta \kappa_{V}=\Delta \kappa_{V}^{\text {upper }}-\Delta \kappa_{V}^{\text {lower }}, \delta \lambda_{V}=\lambda_{V}^{\text {upper }}-\lambda_{V}^{\text {lower }} \tag{4}
\end{equation*}
$$

The current limits on anomalous couplings and the difference of the upper and lower bounds for electron beam energies of $\mathrm{E}_{\mathrm{e}}=80 \mathrm{GeV}$ with integrated luminosities $100 \mathrm{fb}^{-1}$ at FCC-ep with the unpolarized (polarized) electron beam are given in Table III. We have obtained two-parameter limits on $\delta \Delta \kappa_{\gamma}$ and $\delta \lambda_{\gamma}$ which can be compared to the ATLAS and CMS results. However, the current limits on $\delta \lambda_{Z}$ is found to be much more sensitive at the FCC-ep.

TABLE III
THE 95\% C.L. Current Limits on the Anomalous Couplings and the Difference of the Upper and Lower Bounds for Electron Beam
Energy Of $\mathrm{E}_{\mathrm{E}}=80 \mathrm{GEV}$ wITH $\mathrm{L}_{\mathrm{int}}=100 \mathrm{FB}^{-1}$ FOR POLARIZED ELECTRON BEAM

| Pe | $\Delta \kappa \gamma$ | $\delta \Delta \kappa \gamma$ | $\lambda \gamma$ | $\delta \lambda \gamma$ |
| :---: | :---: | :---: | :---: | :---: |
| -0.8 | $-0.100: 1.001$ | 1.101 | $-0.026: 0.039$ | 0.0650 |
| Pe | $\Delta \kappa \mathrm{z}$ | $\delta \Delta \kappa \mathrm{z}$ | $\lambda \mathrm{z}$ | $\delta \lambda \mathrm{z}$ |
| -0.8 | $-0.019: 0.301$ | 0.320 | $-0.0011: 0.0012$ | 0.0023 |



Fig. 13 Two dimensional 95\% C.L contour plot anomalous couplings in the $\lambda_{\gamma}-\Delta \kappa_{\gamma}$ plane for the integrated luminosity of $10 \mathrm{fb}^{-1}$ and 100 $\mathrm{fb}^{-1}$ at electron beam energy $E_{e}=80 \mathrm{GeV}$ with polarization $P_{e}=-0.8$


Fig. 14 Two-dimensional 95\% C.L contour plot of anomalous couplings in the $\lambda_{Z}-\Delta \kappa_{Z}$ plane for the integrated luminosity of $10 \mathrm{fb}^{-1}$ and $100 \mathrm{fb}^{-1}$ at electron beam energy $E_{e}=80 \mathrm{GeV}$ with polarization

## VII. CONCLUSION

The $W W \gamma$ and $W W Z$ anomalous interactions through the processes $e p \rightarrow v e^{q \gamma X}$ and $e p \rightarrow v e^{q Z X}$ can be studied independently at the LHeC and FCC-ep. We obtain twoparameter accessible ranges of triple gauge boson anomalous couplings at LHeC and FCC -ep with the polarized electron beam at the energies $E_{e}=140 \mathrm{GeV}$ and $\mathrm{E}_{\mathrm{p}}=7 \mathrm{TeV}$, and $E_{e}=80$ GeV and $\mathrm{E}_{\mathrm{p}}=50 \mathrm{TeV}$, respectively. Our limits compare with the results from two-parameter analysis given by ATLAS and CMS Collaborations [10]-[13]. We find that the sensitivities to anomalous couplings $\Delta \kappa_{V}(V=\gamma, Z)$ will be of the order of $10^{-1}$, which is an order of magnitude larger than the SM loop level sensitivity of $10^{-2}$, however a measurement of these couplings above $10^{-2}$ would offer a possible new physics signal. We conclude that the anomalous couplings $\lambda_{\gamma}$ and $\lambda_{Z}$ can be well constrained with the sensitivity of the order of $10^{-2}$ and $10^{-3}$ at the FCC-ep with polarized electron beam. The LHeC and FCC-ep could give complementary information about anomalous couplings compared to Tevatron and LHC.

## ACKNOWLEDGMENT

The work of O.C. is partially supported by State Planning Organisation (DPT) - Ministry of Development under the grant No. DPT2006K-120470. A.S. would like to thank Abant Izzet Baysal University Department of Physics where of part this study was carried out for their hospitality

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